

Spatially Covariant Theories of a Transverse, Traceless Graviton

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Cosmic Acceleration



Saul Perlmutter, Brian P. Schmidt, Adam G. Riess

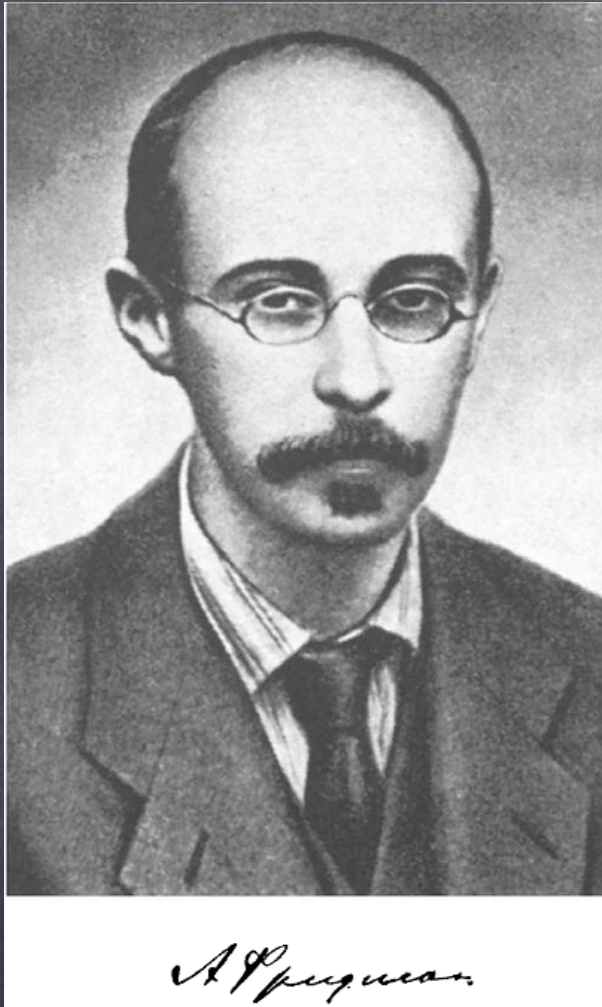
“for the discovery of the accelerating expansion of the Universe through observations of distant supernovae”

$$\ddot{a} > 0$$



2011

Implications of Acceleration



Alexander Friedmann

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3p)$$

$$\ddot{a} > 0 \longrightarrow p < -\frac{1}{3}\rho$$

Cosmological Constant

$$S = \int d^4x \sqrt{-g} \Lambda \longrightarrow p_\Lambda = -\rho_\Lambda$$

Λ CDM, the Concordance Model of Cosmology

$$\Omega_\Lambda \approx 0.7 \longrightarrow \rho_\Lambda \approx (meV)^4$$

Cosmological Constant Problem

Expectation

$$\rho_{\Lambda} \approx M_{Pl}^4$$

Apparent Reality

$$\rho_{\Lambda} \approx (meV)^4 \approx 10^{-120} M_{Pl}^4$$

- **Cosmological Constant Problem** – Why is ρ_{Λ} so small?
 - Cancellation of zero-point energies requires extreme **fine-tuning**
 - Weinberg's **"No-Go"** Theorem: Cannot relax dynamically to $\rho_{\Lambda} \approx 0$

S. Weinberg, Rev.Mod.Phys. 61 (1989) 1-23

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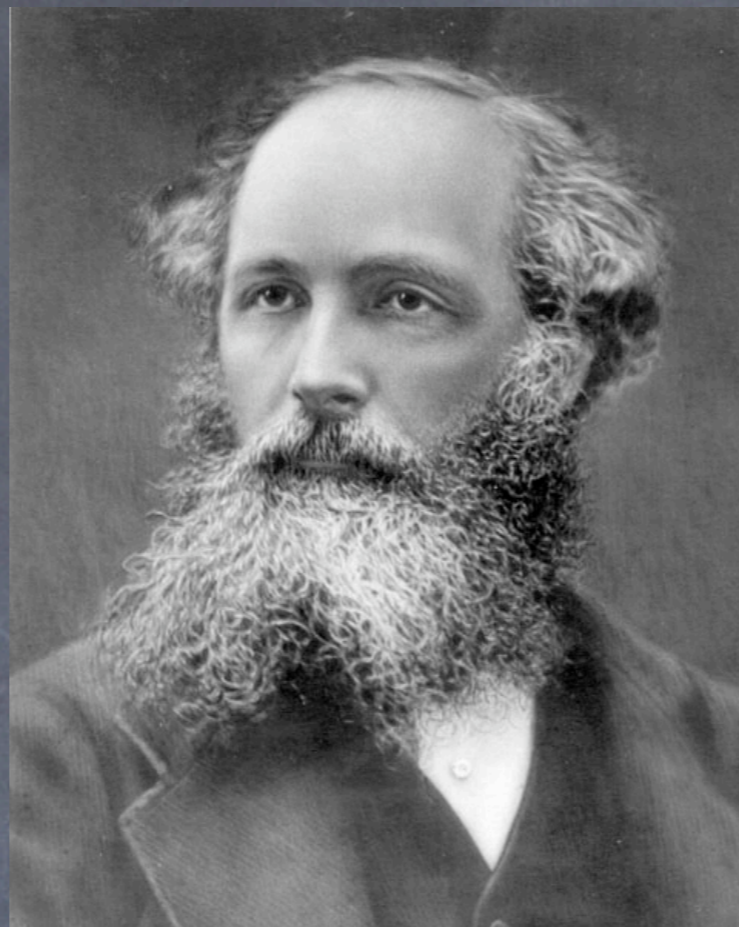
• Possible Solutions

- Dark Energy
- Modified Gravity
- ?

Can General Relativity be modified
without new degrees of freedom?

Degrees of Freedom

- Independent parameters needed to specify the state of a system
- Classical field theories have infinitely many degrees of freedom, but...
- Finitely many **local** degrees of freedom
- After quantization, these count particle **polarization** states



$$S = \int d^4x F_{\mu\nu} F^{\mu\nu}$$

Field

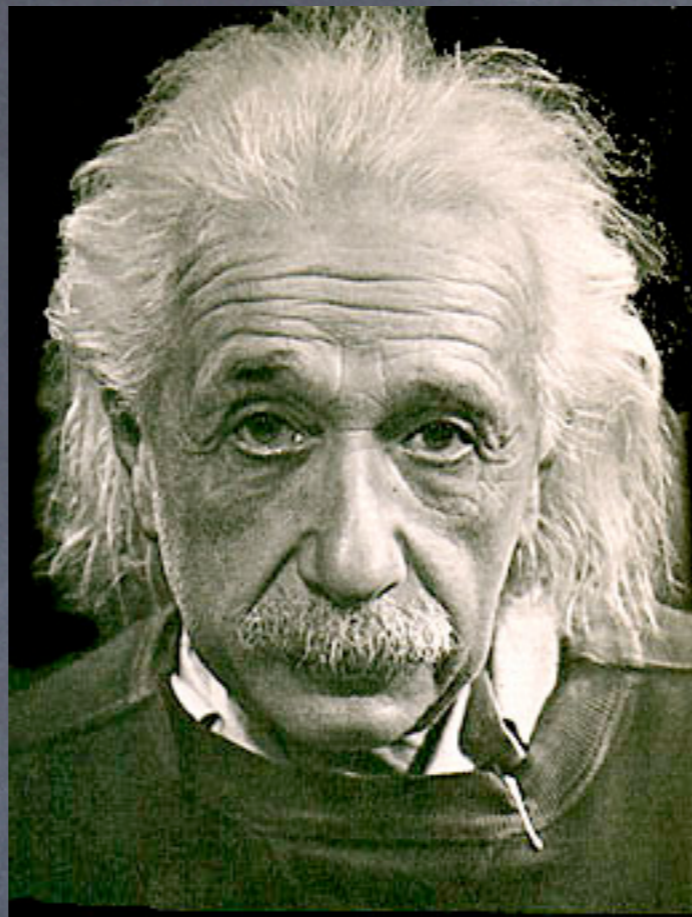
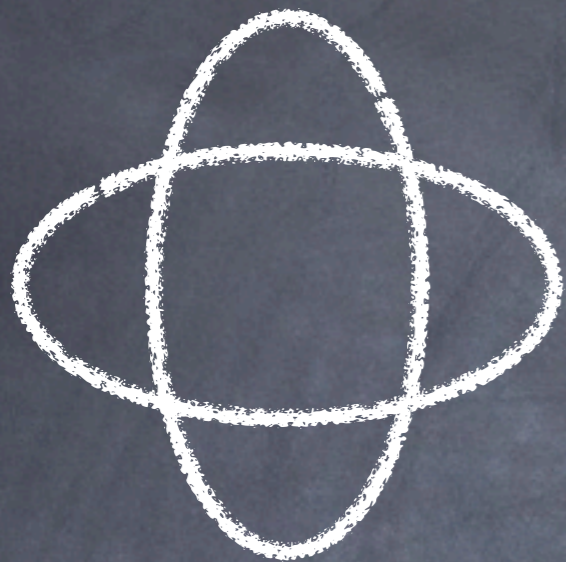
A_μ

Particle

Photon

Polarizations

2



$$S = \int d^4x \sqrt{-g} R$$

Field

$g_{\mu\nu}$

Particle

Graviton

Polarizations

2

Uniqueness Theorem

- In Lorentz covariant theories, polarization states are fixed by **mass** and **spin** of particles
- **Weinberg's Theorem:** GR is the **unique** Lorentz covariant theory of an interacting massless spin-2 particle

S. Weinberg

Phys. Rev. 138 (1965) B988-B1002

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Lorentz covariant modifications of GR
introduce **new degrees of freedom**

Modifying the Graviton

- From binary pulsars, strong evidence for **two** polarizations
- Lorentz covariant modifications of GR introduce **new** particles or polarizations, but we haven't seen any
- To modify the behavior of the graviton **without** new degrees of freedom, we must...

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Break Lorentz covariance...

Modifying the Graviton

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Break Lorentz covariance...

Explicitly!

Cosmic Rest Frame



2006

John C. Mather, George F. Smoot

“for their discovery of the blackbody form and anisotropy of the cosmic microwave background radiation”

Spontaneous or **Explicit** Symmetry Breaking?

ADM Action for GR

$$\begin{array}{ccc} 10 & & 6 \quad 3 \quad 1 \\ g_{\mu\nu} & \longrightarrow & h_{ij} \quad N^i \quad N \end{array}$$

Arnowitt-Deser-Misner Action

$$S = \int dt d^3x N \sqrt{h} \left(K^{ij} K_{ij} - K^2 + R^{(3)} - 2\Lambda \right)$$

Extrinsic Curvature

$$K_{ij} \equiv \frac{1}{2} N^{-1} \left(\dot{h}_{ij} - \nabla_i N_j - \nabla_j N_i \right)$$

→ Only the spatial metric h_{ij} is dynamical

Degrees of Freedom in GR

- 6 dynamical variables h_{ij} constrained by
- 4 spacetime gauge symmetries

$$6 - 4 = 2 \text{ degrees of freedom}$$

- These are the polarizations of the graviton
- To avoid new degrees of freedom, balance dynamical variables against gauge symmetries

Approach

- GR is a theory of a spatial metric h_{ij} (which has **six** components) subject to **four** spacetime gauge symmetries

$$6 - 4 = 2$$

- We consider theories of a **unit-determinant** spatial metric \tilde{h}_{ij} (which has **five** components) subject to **three** spatial gauge symmetries

$$5 - 3 = 2$$

- In a particular gauge, General Relativity can be cast in this form

What freedom is there to modify
spatially covariant General Relativity?

Ultralocal Modified Gravity

- Neglect terms in action of quadratic or higher order in spatial derivatives
- This is a long-wavelength, deep **infrared** limit in which gravitons have only **kinetic** energy
- In this limit, consistency implies only one restriction: **spatial conformal symmetry**

$$x^i \longrightarrow x^i / \lambda$$

Local Modified Gravity

- Allow the action to depend on spatial derivatives through the Ricci scalar \tilde{R} of \tilde{h}_{ij}
- Ricci scalar dependence is the **leading local correction** to infrared dynamics
- This class of theories **includes GR**
- In addition to spatial conformal symmetry, consistency requires the vanishing of a non-trivial field dependent vector quantity

$$\begin{aligned}
 \mathcal{A}_k = & \frac{1}{2} \omega^2 \frac{\partial \pi_\omega}{\partial \tilde{R}} \left(2\tilde{\pi}^{ij} \tilde{\nabla}_i - \omega (\tilde{\nabla}^j \pi_\omega) \right) \left(\sum_{n=2}^{\infty} n \Pi(n-2)_{jk} \frac{\partial \pi_\omega}{\partial \phi(n)} \right) \\
 & - \frac{1}{3} \omega^2 \frac{\partial \pi_\omega}{\partial \tilde{R}} \tilde{\nabla}_k \left(\sum_{n=3}^{\infty} n \phi(n-1) \frac{\partial \pi_\omega}{\partial \phi(n)} \right) - \omega \left(\omega \frac{\partial \pi_\omega}{\partial \tilde{\pi}^{jk}} \tilde{\nabla}^j + \frac{4}{3} \tilde{\nabla}_k \right) \frac{\partial \pi_\omega}{\partial \tilde{R}}
 \end{aligned}$$

Summary of Results

- We **can** modify GR **without** new DOF, provided the theory is invariant under **spatial conformal symmetry**
- In the ultralocal limit, the action is otherwise arbitrary
- Locality constrains the action through the consistency condition

$$\mathcal{A}_k = 0$$

Win-Win Situation

- Is GR the **unique** low-energy **local**, realistic theory of the graviton degrees of freedom?
- If so, Lorentz invariance in the gravitational sector could arise as an **accidental** symmetry
- If not, infrared modifications of the graviton could shed light on **dark energy**
- Either way, the results will be interesting!

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Thank you